

FeynRules Implementation of Z' Model

M.Sc. THESIS

by

Joginder

(Roll No.- 2103151010)



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FeynRules Implementation of Z' Model

A THESIS

*Submitted in partial fulfilment of the
requirements for the award of the degree
of*
Master of Science

by
Joginder



DISCIPLINE OF PHYSICS
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June, 2023



INDIAN INSTITUTE OF TECHNOLOGY
INDORE

CANDIDATE'S DECLARATION

I hereby certify that the work which is being presented in the thesis entitled **FeynRules Implementation of Z' Model** in the partial fulfillment of the requirements for the award of the degree of **Master of Science** and submitted in the **Discipline of Physics, Indian Institute of Technology Indore**, is an authentic record of my own work carried out during the time period from July 2022 to June 2023 under the supervision of **Dr Dipankar Das, Assistant professor, Indian Institute of Technology Indore**.

The matter presented in this thesis has not been submitted by me for the award of any other degree of this or any institute.

Joginder

Signature of the student with date

29/05/23 (Joginder)

This is to certify that the above statement made by the candidate is correct to the best of my knowledge.

Dipankar Das 7/6/23

Signature of the Supervisor of M.Sc. thesis (with date)

(Dr Dipankar Das)

JOGINDER has successfully given his M.Sc. Oral Examination held on

Date

Dipankar Das

Signature of Supervisor of M.Sc. thesis

Date:

7/6/23

Sushanta Ranchar

Signature of PSPC Member no. 1

Date:

7/06/23

Mandendra Nath

Convener, DPGC

Date: 7/6/2023

Sushanta Ranchar

Signature of PSPC Member no. 2

Date:

07/06/23

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Contents

1	Introduction	1
1.1	Implementation Flow	2
2	Z' Model	4
2.1	Gauge Sector	4
2.2	Scalar Sector	7
3	FeynRules	11
3.1	Model Description	12
3.2	Lagrangian	37
3.3	Running The FeynRules	38
4	MadGraph	40
4.1	Installation	40
4.2	Cards Files	45
4.3	Plots	46
5	Results	49
5.1	Future	53

List of Figures

1.1	Flow Chart	2
3.1	Model file describing general model information, vacuum expectation values and Gauge groups (Screenshot 1)	13
3.2	Model file describing indices , interaction order hierarchy and calling of parameter and particle files (Screenshot 2)	14
3.3	Parameter file describing parameters aEWM1, Gf, aS, ymdo, ymup, yms and ymc (Screenshot 1)	16
3.4	Parameter file describing ymb, ymt, yme, ymm, ymtau, cabi and saz (Screenshot 2)	17
3.5	Parameter file describing rs, st, gs, aEW and ee (Screenshot 3)	19
3.6	Parameter file describing vev, Ge2, th, cw, sw, tw and MW (Screenshot 4)	21
3.7	Parameter file describing gw, g1, vevs, dm1, dm2, xpp2 and xpp (Screenshot 5)	24
3.8	Parameter file describing tx, gxp, chi, sc, tc and gx (Screenshot 6)	25
3.9	Parameter file describing ls, muS, lHS, lh, muH and yl (Screenshot 7)	28
3.10	Parameter file describing yu, yd and CKM (Screenshot 8)	29
3.11	Particle file describing physical vector fields A, Z, W and G (Screenshot 1)	30
3.12	Particle file describing fields Zp, B, Bp, X and Xp (Screenshot 2)	31

3.13	Particle file describing fields Wi, vl and l (Screenshot 3)	32
3.14	Particle file describing fields uq, dq, LL and lR (Screenshot 4)	34
3.15	Particle file describing fields QL, uR, dR, H and Hp (Screenshot 5)	35
3.16	Particle file describing fields hd, G0, GP and GZp (Screenshot 6)	36
3.17	Particle file describing fields ss, Phi and Su (Screenshot 7)	37
3.18	Lagrangian	38
4.1	Installation	41
4.2	Feynman Diagram (Screenshot 1)	42
4.3	Feynman Diagram (Screenshot 2)	43
4.4	Feynman Diagram (Screenshot 3)	43
4.5	Feynman Diagram (Screenshot 4)	44
4.6	Event Generation	44
4.7	Sample Run	45
4.8	Script File	47
4.9	Parameter Card	48
5.1	Model Data Plot	49
5.2	Model Data Set	50
5.3	Experimental Data Plot 1	51
5.4	Experimental Data Plot 2	51
5.5	Combined Plot	52
5.6	Model Data Set 2	52
5.7	Final Plot	53

Chapter 1

Introduction

Standard Model has been proved experimentally to be a good model for explaining the behavior of subatomic particles. The gauge group of the Standard Model is $SU(3)_C \times SU(2)_L \times U(1)_Y$. The Standard Model describes three of the four fundamental forces of nature: the strong force, the weak force, and the electromagnetic force. The fundamental particles in the Standard Model are divided into two groups: fermions and bosons. Fermions include quarks and leptons, which make up all matter. Bosons include force carriers such as photons, W and Z bosons, and gluons. In the scalar sector, Standard Model have Higgs Doublet. But there are some phenomena which cannot be explained on the basis of Standard Model. There are many theories and models trying to explain these phenomena. One of these is the Z' model which is an extension of the Standard Model of particle physics. Here, I have worked on a minimal Z' model[1]. The gauge group for Z' model is $SU(3)_C \times SU(2)_L \times U(1)_Y \times U(1)_X$. There is an extra gauge boson in Z' model in addition to the ones in Standard Model. In the scalar sector, Z' Model (used here) have Higgs Doublet and an additional Higgs scalar. This also adds an extra Higgs particle apart from that of the Standard Model. The model can be implemented using FeynRules which is a powerful tool for generating Feynman rules and amplitudes for particle physics models. Feynrules requires model file and Lagrangian written in the Wolfram language. Model file contains various classes like particle class, parameter class, Gauge group class etc. We can

have a single model file or different files separately like parameter file and particle file. The Lagrangian is written during the mathematica session. The output of FeynRules is input to MadGraph which is a software package to generate simulated events. Events generated at colliders can be simulated with the help of Monte Carlo generators. Madgraph can be used to obtain cross-section and various histogram plots. By selecting a suitable process involving Z' particle, we can obtain the cross-section results. If mass of the Z' particle is changed, the cross-section also changes. We can draw graph of cross-section versus mass of the Z' particle. This graph from simulation can be compared with the graph obtained from experiments[2].

1.1 Implementation Flow

The implementation flow of the procedure used here can be understood from the following diagram :-

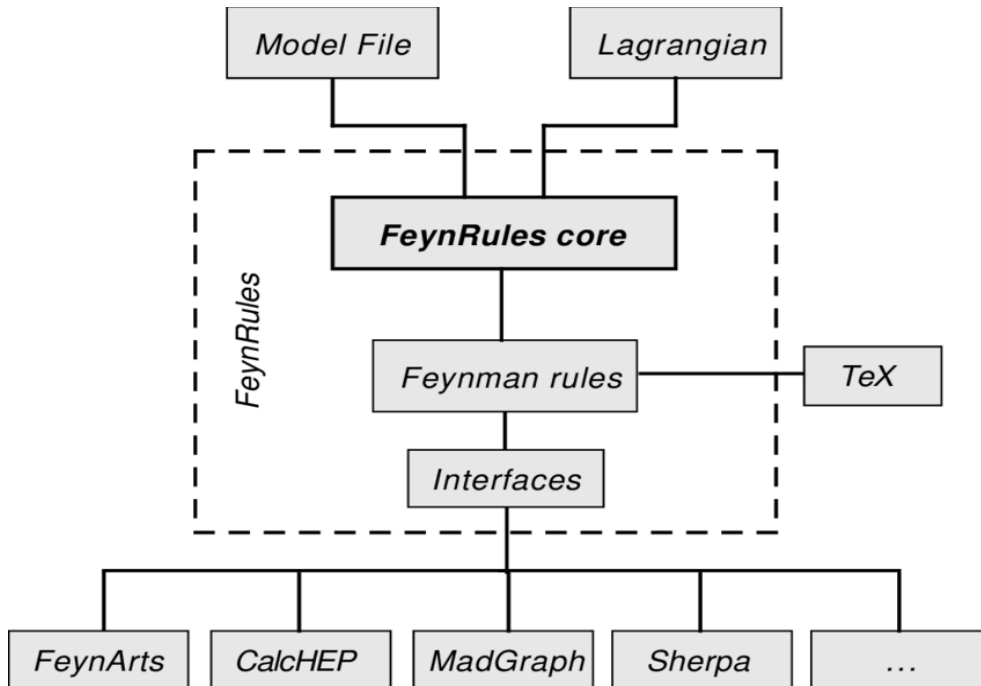


Figure 1.1: Flow Chart

FeynRules is a Mathematica-based package which addresses the implementation of particle physics models. For implementing a new model in FeynRules, we need two things : Model File and Lagrangian. In Feyn-

Rules, a model file is a text file that contains the information needed to define a particular model, including its particle content, interactions, and symmetry properties. The Lagrangian is a mathematical function that describes the dynamics of a system. It includes the interaction terms which define a particular theory. The model file have various classes. The classes have information of fields, parameters and gauge group of model which. All the fields including scalar, Dirac and vector fields are defined in a particle class. After running the FeynRules, a FeynRule file is created. It is written in the Universal FeynRules Output (UFO) format (interface), which is a standardized format used by various software tools to generate Feynman diagrams and perform other calculations related to the model. The FeynRules file typically includes information about the particle content of the model, including the masses, quantum numbers, and interactions of each particle. It also includes information about the vertices and couplings between particles, as well as any additional parameters that are relevant for the model. A Feynrule can be now used for anyone of the Monte Carlo generators. We don't have to write different codes for different Monte Carlo generators.

Chapter 2

Z' Model

In Beyond Standard Model physics, Z' Models are the most common. Here, we will discuss minimal Z' model. We have to add one more symmetry group ($U(1)$) which is the simplest thing we can do. The gauge group is $SU(3)_C \times SU(2)_L \times U(1)_Y \times U(1)_X$. The consequence is that we get a new vector boson Z' in addition to three vector bosons of the standard model.

2.1 Gauge Sector

Since we have an extra symmetry group in Z' Model, there are four gauge bosons in the theory. The Lagrangian[3] is given by :-

$$L_{Gauge} = -\frac{1}{4}W_a^{\mu\nu}W_{\mu\nu}^a - \frac{1}{4}B^{\mu\nu}B_{\mu\nu} - \frac{1}{4}X^{\mu\nu}X_{\mu\nu} - \frac{\sin\chi}{2}B^{\mu\nu}X_{\mu\nu} \quad (2.1)$$

Here, $W_a^{\mu\nu}$ is for $SU(2)$ group (a goes from 1 to 3 for three W gauge bosons), $B^{\mu\nu}$ for $U(1)_Y$ and $X^{\mu\nu}$ for $U(1)_X$ symmetry group. The last term in Lagrangian depicts that we have a gauge kinetic mixing term for two $U(1)$ group gauge bosons and it can be controlled through the χ parameter. This gauge kinetic mixing term is the only possible gauge invariant term for the Lagrangian. So, all the other possible combinations are not allowed. But, we want our kinetic terms in the Lagrangian to be canonically diagonal. So, we perform a linear transformation as following :-

$$\begin{pmatrix} B'_\mu \\ X'_\mu \end{pmatrix} = \begin{pmatrix} 1 & \sin \chi \\ 0 & \cos \chi \end{pmatrix} \begin{pmatrix} B_\mu \\ X_\mu \end{pmatrix} \quad (2.2)$$

Such a transformation leads to a U(1) charge shift[4]. The inverse transformation can be written as follows :-

$$B_\mu = B'_\mu - \tan \chi X'_\mu \quad (2.3)$$

$$X_\mu = \sec \chi X'_\mu \quad (2.4)$$

The definitions for covariant derivatives are as :-

$$D^\mu \Phi = (\partial_\mu - ig \frac{\tau_a}{2} W_\mu^a - i \frac{g_Y}{2} B_\mu - ig_x \frac{x_\phi}{2} X_\mu) \Phi \quad (2.5)$$

$$D^\mu S = (\partial_\mu - i \frac{g_x}{2} X_\mu) S \quad (2.6)$$

After the linear transformation, the expressions for covariant derivatives are changed which are as following :-

$$D^\mu \Phi = (\partial_\mu - i \frac{g}{2} (\tau_a W_\mu^a - \tan \theta_w B'_\mu - \tan \theta_x x'_\phi X'_\mu)) \Phi \quad (2.7)$$

$$D^\mu S = (\partial_\mu - ig \tan \theta_x \frac{x'_\phi}{2} X'_\mu) S \quad (2.8)$$

Where we have defined,

$$\tan \theta_w = \frac{g_Y}{g} \quad (2.9)$$

$$\tan \theta_x = \frac{g'_x}{g} \quad (2.10)$$

$$g'_x = g_x \sec \chi \quad (2.11)$$

$$x'_\phi = x_\phi - \frac{g_Y}{g_x} \sin \chi \quad (2.12)$$

After spontaneous symmetry breaking, the scalar fields in the unitary gauge can be written as following:-

$$\Phi = \begin{pmatrix} 0 \\ \frac{v+\phi_0}{\sqrt{2}} \end{pmatrix} \quad (2.13)$$

$$S = \frac{v_s + s}{\sqrt{2}} \quad (2.14)$$

where v and v_s are vacuum expectation values of Φ and S respectively. The mass terms for the gauge bosons can be found out as following :-

$$L = (D^\mu \Phi)^\dagger (D_\mu \Phi) + (D^\mu S)^\dagger (D_\mu S) \quad (2.15)$$

We can write the neutral gauge boson mass matrix in the basis where the gauge kinetic terms are diagonal (means selecting W_μ^3, B'_μ and X'_μ). The next step is to make this matrix diagonal. Here, unphysical gauge bosons W_μ^3, B'_μ and X'_μ further mix to give the required physical gauge bosons A_μ, Z_μ and Z'_μ . Since there are three fields to be changed, we need two rotation angles here (namely Weinberg angle(θ_w) and Z-Z' mixing angle(α_z)).

$$L_{Mass} = \frac{1}{2} \begin{pmatrix} W_\mu^3 & B'_\mu & X'_\mu \end{pmatrix} M^2 \begin{pmatrix} W_\mu^3 \\ B'_\mu \\ X'_\mu \end{pmatrix} \quad (2.16)$$

where M^2 is given by :-

$$M^2 = \frac{g^2 v^2}{4} \begin{pmatrix} 1 & -\tan(\theta_w) & -x'_\phi \tan(\theta_x) \\ -\tan(\theta_w) & \tan^2(\theta_w) & x'_\phi \tan(\theta_x) \tan(\theta_w) \\ -x'_\phi \tan(\theta_x) & x'_\phi \tan(\theta_x) \tan(\theta_w) & \tan^2(\theta_x)(r^2 + x'^2_\phi) \end{pmatrix} \quad (2.17)$$

The diagonalization can be done as following :-

$$M^2 = O_g^T M_D^2 O_g \quad (2.18)$$

Where M_D^2 is the squared diagonal gauge boson mass matrix. The diagonalization is a two step process here with rotation angles namely Weinberg angle(θ_w) and Z-Z' mixing angle(α_z). Here, M_D^2 and O_g are given by :-

$$M_D^2 = \begin{pmatrix} M_A^2 & 0 & 0 \\ 0 & M_Z^2 & 0 \\ 0 & 0 & M_{Z'}^2 \end{pmatrix} \quad (2.19)$$

$$O_g = \begin{pmatrix} \cos \theta_w & \sin \theta_w & 0 \\ -\sin \theta_w \sin \alpha_z & \cos \theta_w \cos \alpha_z & \sin \alpha_z \\ \sin \theta_w \sin \alpha_z & -\cos \theta_w \sin \alpha_z & \cos \alpha_z \end{pmatrix} \quad (2.20)$$

Photon mass (M_A) is zero in the equation 2.19. The expression for the transformed fields are given by :-

$$B'_\mu = \cos \theta_w A_\mu - \sin \theta_w \cos \alpha_z Z_\mu + \sin \theta_w \sin \alpha_z Z'_\mu \quad (2.21)$$

$$W_\mu^3 = \sin \theta_w A_\mu + \cos \theta_w \cos \alpha_z Z_\mu - \cos \theta_w \sin \alpha_z Z'_\mu \quad (2.22)$$

$$X'_\mu = \sin \alpha_z Z_\mu + \cos \alpha_z Z'_\mu \quad (2.23)$$

2.2 Scalar Sector

In the standard model, we have a scalar doublet (Higgs Doublet). Here in the Z' Model, we have an additional singlet scalar. The $U(1)_X$ group is broken by this singlet scalar. The scalars are defined as following :-

$$\phi = (1, 2, 1/2, \frac{x_\phi}{2}) \quad (2.24)$$

$$S = (1, 1, 0, 1/2) \quad (2.25)$$

The quantities inside the parentheses characterize the transformation properties under the gauge group of the Z' Model. The first two numbers tell that whether the multiplet is a singlet or a doublet under $SU(3)_C$ and $SU(2)_L$ groups. The last two numbers are hypercharges of $U(1)_Y$ and $U(1)_X$ groups respectively. The Lagrangian for the scalar sector is given by :-

$$L_{Scalar} = (D^\mu \Phi)^\dagger (D_\mu \Phi) + (\overline{D^\mu S})^\dagger (D_\mu S) - V(\Phi, S) \quad (2.26)$$

Where $V(\Phi, S)$, Φ and S are given by :-

$$V(\Phi, S) = -\mu_\phi^2 (\Phi^\dagger \Phi) - \mu_s^2 (S^\dagger S) + \lambda_\phi (\Phi^\dagger \Phi)^2 + \lambda_s (S^\dagger S)^2 + \lambda_{\phi s} (\Phi^\dagger \Phi) (S^\dagger S) \quad (2.27)$$

$$\Phi = \begin{pmatrix} -iG_W \\ \frac{v+\phi_0+iG_Z}{\sqrt{2}} \end{pmatrix} \quad (2.28)$$

$$S = \frac{v_s + s + iG_{Z'}}{\sqrt{2}} \quad (2.29)$$

Here, G_θ , G_P and $G_{Z'}$ are the Goldstone bosons corresponding to Z , W and Z' bosons respectively. In the Unitary Gauge and after spontaneous symmetry breaking, Φ and S become as following :-

$$\Phi = \begin{pmatrix} 0 \\ \frac{v+\phi_0}{\sqrt{2}} \end{pmatrix} \quad (2.30)$$

$$S = \frac{v_s + s}{\sqrt{2}} \quad (2.31)$$

When the equations 2.30 and 2.31 are put in equation 2.27, the mass terms have mixing term for fields ϕ_0 and s . So, we need to diagonalize the mass matrix to get the physical Higgs Scalars. The mass terms in the Lagrangian are as following :-

$$V_{Scalar} = \lambda_\phi v^2 \times \phi_0^2 + \lambda_{\phi s} v \times v_s \times s \times \phi_0 + \frac{\lambda_s v_s^2 \times s^2}{2} \quad (2.32)$$

The diagonalization involves rotation through one angle only (namely θ_s) which is done as following :-

$$\begin{pmatrix} 2 \times \lambda_\phi v^2 & \lambda_{\phi s} v \times v_s \\ \lambda_{\phi s} v \times v_s & 2 \times \lambda_s v_s^2 \end{pmatrix} = O_s^T \begin{pmatrix} M_H^2 & 0 \\ 0 & M_{H'}^2 \end{pmatrix} O_s \quad (2.33)$$

Here, M_H and $M_{H'}$ are the masses of physical Higgs masses. O_s is the rotation matrix with Higgs mixing angle (θ_s) defined as following :-

$$O_s = \begin{pmatrix} \cos \theta_s & \sin \theta_s \\ -\sin \theta_s & \cos \theta_s \end{pmatrix} \quad (2.34)$$

From the equations 2.33 and 2.34, we get

$$\begin{pmatrix} 2 \times \lambda_\phi v^2 & \lambda_{\phi s} v \times v_s \\ \lambda_{\phi s} v \times v_s & 2 \times \lambda_s v_s^2 \end{pmatrix} = \begin{pmatrix} M_H^2 \cos^2 \theta_s + M_{H'}^2 \sin^2 \theta_s & \sin \theta_s \cos \theta_s (M_H^2 - M_{H'}^2) \\ \sin \theta_s \cos \theta_s (M_H^2 - M_{H'}^2) & M_H^2 \sin^2 \theta_s + M_{H'}^2 \cos^2 \theta_s \end{pmatrix} \quad (2.35)$$

The coefficients of fields in the scalar potential (in equation 2.32) are related to the physical Higgs masses. The relations are as following :-

$$2 \times \lambda_\phi v^2 = M_H^2 \cos^2 \theta_s + M_{H'}^2 \sin^2 \theta_s \quad (2.36)$$

$$2 \times \lambda_s v_s^2 = M_H^2 \sin^2 \theta_s + M_{H'}^2 \cos^2 \theta_s \quad (2.37)$$

$$\lambda_{\phi s} v \times v_s = \sin \theta_s \cos \theta_s (M_H^2 - M_{H'}^2) \quad (2.38)$$

The Spontaneous Symmetric Breaking occur as following :-

$$SU(2)_L \times U(1)_Y \times U(1)_X \rightarrow U(1)_{EM} \quad (2.39)$$

Here we can see that only the electromagnetic symmetry group is left unbroken. This is evident from the selection of electric charge expression as

following :-

$$Q = T_{3L} + Y \tag{2.40}$$

, where T_{3L} and Y are the third generator of $SU(2)_L$ and the hypercharge related to $U(1)_Y$ respectively.

Chapter 3

FeynRules

FeynRules is a Mathematica-based package which addresses the implementation of particle physics models and provides a set of functions for defining particle content, interactions, and parameters of a new model. We have used the **Wolfram Mathematica 12.0** version here. It calculates the underlying Feynman rules and outputs them to a form appropriate for various programs such as CalcHep, FeynArts, MadGraph5, Sherpa and Whizard. Since Feynman rules produced can be output to various event generators, an interface is required (e.g. UFO Interface). Here, I have implemented Z' Model. We need FeynRules for the following reasons:

1. Each event generator uses its own syntax for a new model. Learning one syntax did not transfer to another generator.
2. Implementing a new model often required the modification of the event generator code itself. These implementations did not transfer well between theorists or to experimentalists.

FeynRules Specifications

Version: 2.3.49

Authors: A. Alloul, N. Christensen, C. Degrande, C. Duhr, B. Fuks

Website: <http://feynrules.phys.ucl.ac.be>

3.1 Model Description

To use FeynRules [5], the user must start by entering the details of the new model in a valid mathematica syntax (Wolfram Language). Model file contains all of the information needed for a model or a theory. These definitions can be placed in a pure text file with **.fr** extension. The different type of information is included in various classes. All classes can be in a single file or we can make separate files for classes (e.g. for parameters, fields(particles), etc.). The main model file **ZPrime.fr** is shown in figure 3.1 and 3.2 .

```

1  (***** This is the FeynRules mod-file for Z-prime Model *****)
2  (***** Author: Joginder *****)
3
4
5  M$ModelName = "Z Prime Model";
6
7  (***** Model Information *****)
8
9  M$Information = {
10   Authors    -> {"Joginder"},
11   Version    -> "1.0",
12   Date       -> "01. 01. 2023",
13   Institutions -> "Indian Institute of Technology Indore",
14   Emails     -> {"jogindergoswami77@gmail.com"}
15 };
16
17 FeynmanGauge = True;
18
19 (* ***** NLO Variables ***** *)
20
21 FR$LoopSwitches = {{Gf, MW}};
22 FR$RmDbExt = { ymb -> MB, ymc -> MC, ymdo -> MD, yme -> Me,
23   ymm -> MMU, yms -> MS, ymt -> MT, ymtau -> MTA, ymup -> MU};
24
25 (* ***** vevs ***** *)
26
27 M$vevs = { {Phi[2],vev},{Su,vevs} };
28
29 (* ***** Gauge groups ***** *)
30
31 M$GaugeGroups = {
32   U1Y == {
33     Abelian          -> True,
34     CouplingConstant -> g1,
35     GaugeBoson       -> B,
36     Charge           -> Y
37   },
38   (* =====New Group U1X Added===== *)
39   U1X == {
40     Abelian          -> True,
41     CouplingConstant -> gx,
42     GaugeBoson       -> X,
43     Charge           -> Yp
44   },
45   SU2L == {
46     Abelian          -> False,
47     CouplingConstant -> gw,
48     GaugeBoson       -> Wi,
49     StructureConstant -> Eps,
50     Representations  -> {Ta,SU2D},
51     Definitions      -> {Ta[a_,b_,c_] -> PauliSigma[a,b,c]/2, FSU2L[i_,j_,k_] -> I Eps[i,j,k]}
52   },
53   SU3C == {
54     Abelian          -> False,
55     CouplingConstant -> gs,
56     GaugeBoson       -> G,
57     StructureConstant -> f,
58     Representations  -> {T,Colour},
59     SymmetricTensor  -> dSUN
60   }
61 };

```

Figure 3.1: Model file describing general model information, vacuum expectation values and Gauge groups (Screenshot 1)

```

63  (* ***** Indices ***** *)
64
65  IndexRange[Index[SU2W      ]] = Unfold[Range[3]];
66  IndexRange[Index[SU2D      ]] = Unfold[Range[2]];
67  IndexRange[Index[Gluon     ]] = NoUnfold[Range[8]];
68  IndexRange[Index[Colour     ]] = NoUnfold[Range[3]];
69  IndexRange[Index[Generation]] = Range[3];
70
71  IndexStyle[SU2W,      j];
72  IndexStyle[SU2D,      k];
73  IndexStyle[Gluon,     a];
74  IndexStyle[Colour,    m];
75  IndexStyle[Generation, f];
76
77  (* *** Interaction orders *** *)
78  (* *** (as used by mg5) *** *)
79
80  M$InteractionOrderHierarchy = {
81  {QCD, 1},
82  {QED, 2}
83  };
84
85  (***** Calling for paramter and particles files ***** )
86
87  Get["parameterszp.fr"];
88
89  Get["particleszp.fr"];

```

Figure 3.2: Model file describing indices , interaction order hierarchy and calling of parameter and particle files (Screenshot 2)

1. Model Information

Just like any programming language, the model information is stored in some variables. The variables in Wolfram language are defined with **M\$Variablename** This contains general information about the model. It includes model name. The model information have options such as authors, institute, date and email IDs etc. The variable **M\$Information** acts as an electronic signature of the model. **Lorentz and Spin** indices are defined by default in FeynRules. Indices are defined by **Index[]**.

2. Gauge Group

The structure of the interactions of a model is in general dictated by gauge symmetries. In Z' model, Gauge group is $SU(3) \times SU(2) \times U(1) \times U(1)$. The options used in this class are Abelian, Coupling Constant, Charge etc. Since, we have doublets for $SU(2)$ group, generators are used in the two-dimensional representation $(\frac{\tau_a}{2})$. These are defined in line number 51.

3. Model Parameters

Lagrangian is in terms of parameters and fields. So, for listing all the parameters, we need a parameter class. Parameters can be internal or external. The parameters directly taken from the experiments are external and the parameters derived from the external or/and internal parameters are internal parameters. There are various options for a particular parameter like value, parameter type, complex parameter type, interaction order etc. There are total 46 parameters in Z' model. The external parameters are defined from line numbers **7** to **150**. Apart from the Standard Model model file, we have four new external parameters in Z' model namely sa_z , r_s , st_s and tx . Here, the details are included in **parameterszp.fr** file which is given by :-

```

1  (* ***** Parameters ***** *)
2
3  M$Parameters = {
4
5      (* External parameters *)
6
7      aEWM1 == {
8          ParameterType -> External,
9          BlockName     -> SMINPUTS,
10         OrderBlock    -> 1,
11         Value          -> 127.9,
12         InteractionOrder -> {QED,-2},
13         Description    -> "Inverse of the EW coupling constant at the Z pole"
14     },
15     Gf == {
16         ParameterType -> External,
17         BlockName     -> SMINPUTS,
18         OrderBlock    -> 2,
19         Value          -> 1.16637*^-5,
20         InteractionOrder -> {QED,2},
21         TeX            -> Subscript[G,f],
22         Description    -> "Fermi constant"
23     },
24     aS == {
25         ParameterType -> External,
26         BlockName     -> SMINPUTS,
27         OrderBlock    -> 3,
28         Value          -> 0.13,
29         InteractionOrder -> {QCD,2},
30         TeX            -> Subscript[\[Alpha],s],
31         Description    -> "Strong coupling constant at the Z pole"
32     },
33     ymdo == {
34         ParameterType -> External,
35         BlockName     -> YUKAWA,
36         OrderBlock    -> 1,
37         Value          -> 5.04*^-3,
38         Description    -> "Down Yukawa mass"
39     },
40     ymup == {
41         ParameterType -> External,
42         BlockName     -> YUKAWA,
43         OrderBlock    -> 2,
44         Value          -> 2.55*^-3,
45         Description    -> "Up Yukawa mass"
46     },
47     yms == {
48         ParameterType -> External,
49         BlockName     -> YUKAWA,
50         OrderBlock    -> 3,
51         Value          -> 0.101,
52         Description    -> "Strange Yukawa mass"
53     },
54     ymc == {
55         ParameterType -> External,
56         BlockName     -> YUKAWA,
57         OrderBlock    -> 4,
58         Value          -> 1.27,
59         Description    -> "Charm Yukawa mass"
60     },

```

Figure 3.3: Parameter file describing parameters aEWM1, Gf, aS, ymdo, ymup, yms and ymc (Screenshot 1)

```

61 ymb == {
62   ParameterType -> External,
63   BlockName     -> YUKAWA,
64   OrderBlock    -> 5,
65   Value         -> 4.7,
66   Description    -> "Bottom Yukawa mass"
67 },
68 ymt == {
69   ParameterType -> External,
70   BlockName     -> YUKAWA,
71   OrderBlock    -> 6,
72   Value         -> 172,
73   Description    -> "Top Yukawa mass"
74 },
75 yme == {
76   ParameterType -> External,
77   BlockName     -> YUKAWA,
78   OrderBlock    -> 11,
79   Value         -> 5.11*^-4,
80   Description    -> "Electron Yukawa mass"
81 },
82 ymm == {
83   ParameterType -> External,
84   BlockName     -> YUKAWA,
85   OrderBlock    -> 13,
86   Value         -> 0.10566,
87   Description    -> "Muon Yukawa mass"
88 },
89 ymtau == {
90   ParameterType -> External,
91   BlockName     -> YUKAWA,
92   OrderBlock    -> 15,
93   Value         -> 1.777,
94   Description    -> "Tau Yukawa mass"
95 },
96 cabi == {
97   ParameterType -> External,
98   BlockName     -> CKMBLOCK,
99   OrderBlock    -> 1,
100  Value         -> 0.227736,
101  TeX           -> Subscript[\[Theta], c],
102  Description    -> "Cabibbo angle"
103 },
104
105 (*=====saz Added=====*)
106
107 saz == {
108   ParameterType -> External,
109   BlockName     -> MIXBLOCK,
110   OrderBlock    -> 1,
111   Value         -> 0.0001,
112   TeX           -> Subscript[sa,z],
113   Description    -> "Sine of Z-Zp Mixing Angle"
114 },

```

Figure 3.4: Parameter file describing ymb, ymt, yme, ymm, ymtau, cabi and saz (Screenshot 2)

Here, g_s , α_{EW} and ee are strong coupling constant in natural units, fine structure constant in natural units and electric coupling constant respectively. ' ee ' defined here is equivalent to electric charge(e). ' α_s ' is the strong coupling constant at the Z pole (an experimental or external parameter). At the line numbers **154**, **163** and **170**, g_s , α_{EW} and ee are defined respectively as :-

$$g_s = \sqrt{4\pi \times \alpha_s} \quad (3.1)$$

$$\alpha_{EW} = \frac{1}{a_{EW}M1} \quad (3.2)$$

$$ee = \sqrt{4\pi \times \alpha_{EW}} \quad (3.3)$$

```

117 (*=====rs added=====*)
118
119 rs == {
120   ParameterType -> External,
121   BlockName     -> MIXBLOCK,
122   OrderBlock    -> 2,
123   Value         -> 500,
124   TeX           -> Subscript[r,s],
125   Description    -> "Ratio of Singlet to Doublet VEVs"
126 },
127
128 (*=====st added=====*)
129
130 st == {
131   ParameterType -> External,
132   BlockName     -> MIXBLOCK,
133   OrderBlock    -> 6,
134   Value         -> 0.2,
135   TeX           -> Subscript[st,s],
136   Description    -> "Sin of Higgs Mixing Angle"
137 },
138
139 (*=====tx Added=====*)
140
141 tx == {
142   ParameterType -> Internal,
143   BlockName     -> MIXBLOCK,
144   OrderBlock    -> 7,
145   Value         -> 1,
146   TeX           -> tx,
147   Description    -> "Tan of U1X Mixing Angle"
148 },
149
150
151 (* Internal Parameters *)
152 gs == {
153   ParameterType -> Internal,
154   Value         -> Sqrt[4 Pi aS],
155   InteractionOrder -> {QCD,1},
156   TeX           -> Subscript[g,s],
157   ParameterName  -> G,
158   Description    -> "Strong coupling constant at the Z pole"
159 },
160
161 aEW == {
162   ParameterType -> Internal,
163   Value         -> 1/aEWM1,
164   InteractionOrder -> {QED,2},
165   TeX           -> Subscript[\[Alpha], EW],
166   Description    -> "Electroweak coupling constant"
167 },
168 ee == {
169   ParameterType -> Internal,
170   Value         -> Sqrt[4 Pi aEW],
171   InteractionOrder -> {QED,1},
172   TeX           -> e,
173   Description    -> "Electric coupling constant"
174 },
175

```

Figure 3.5: Parameter file describing rs, st, gs, aEW and ee (Screenshot 3)

Fermi constant (G_f) is related to vacuum expectation value (vev) of scalar doublet and vev is expressed in terms of G_f at line number 180.

$$G_f = \frac{1}{\sqrt{2} \times vev^2} \quad (3.4)$$

We know, Weinberg angle(θ_w) is related to mass of W and Z bosons as :

$$\cos \theta_w = \frac{M_W}{M_Z} \quad (3.5)$$

$$M_W = \frac{g_w \times vev}{2} \quad (3.6)$$

' g'_w is weak coupling constant of $SU(2)_L$ group (At the line number 47 in ZPrime.fr). From equations 3.5 and 3.6, $Ge2$ is defined at line number **188** as :-

$$Ge2 = \frac{g_w}{2 \times \cos g_w} = \frac{M_Z}{vev} \quad (3.7)$$

We know,

$$ee = g_w \times \sin \theta_w \quad (3.8)$$

Using equations 3.5, 3.7 and 3.8, we get

$$\sin 2\theta_w = \frac{ee}{Ge2} \quad (3.9)$$

From equation 3.9, Weinberg angle (θ_w) is defined at line number **196** (th) as :-

$$\theta_w = \frac{\arcsin \frac{ee}{Ge2}}{2} \quad (3.10)$$

From th (θ_w), internal parameters c_w , s_w and t_w are defined at line numbers **204**, **212** and **220** respectively. Parameter M_W is defined at the line number **229** as :-

$$M_W = M_Z \times c_w \quad (3.11)$$

```

177 (*=====vev modified=====*)
178 vev == {
179   ParameterType -> Internal,
180   Value         -> Sqrt[1/(Sqrt[2]*Gf)],
181   InteractionOrder -> {QED,-1},
182   Description    -> "Scalar Doublet vacuum expectation value"
183 },
184
185 (*=====Ge2 added=====*)
186 Ge2 == {
187   ParameterType -> Internal,
188   Value         -> MZ/vev,
189   InteractionOrder -> {QED,1},
190   Description    -> "Extra paramter 1"
191 },
192
193 (*=====th added=====*)
194 th == {
195   ParameterType -> Internal,
196   Value         -> ArcSin[ee/Ge2]/2,
197   TeX           -> Subscript[\[Theta], w],
198   Description    -> "Weinberg angle"
199 },
200
201 (*=====cw modified=====*)
202 cw == {
203   ParameterType -> Internal,
204   Value         -> Cos[th],
205   TeX           -> Subscript[c,w],
206   Description    -> "Cosine of the Weinberg angle"
207 },
208
209 (*=====sw modified=====*)
210 sw == {
211   ParameterType -> Internal,
212   Value         -> Sin[th],
213   TeX           -> Subscript[s,w],
214   Description    -> "Sine of the Weinberg angle"
215 },
216
217 (*=====tw added=====*)
218 tw == {
219   ParameterType -> Internal,
220   Value         -> Tan[th],
221   TeX           -> Subscript[t,w],
222   Description    -> "Tan of the Weinberg angle"
223 },
224
225 (*=====W Mass Modified=====*)
226 MW == {
227   ParameterType -> Internal,
228   Value         -> MZ*cw,
229   TeX           -> Subscript[M,W],
230   Description    -> "W mass"
231 },
232

```

Figure 3.6: Parameter file describing vev, Ge2, th, cw, sw, tw and MW (Screenshot 4)

We know, mass of W boson is given by :-

$$M_W = \frac{g_w \times vev}{2} \quad (3.12)$$

From equations 3.11 and 3.12, g_w is defined at the line number **237** as:-

$$g_w = \frac{2 \times MW}{vev} \quad (3.13)$$

$U(1)_Y$ coupling constant (g_1) is the coupling constant of $U(1)_Y$ group (At the line number 34 in ZPrime.fr). $U(1)_Y$ coupling constant (g_1) and scalar singlet vacuum expectation value ($vevs$) are defined at line numbers **247** and **256** respectively as :-

$$g_1 = g_w \times t_w \quad (3.14)$$

$$vevs = r_s \times vev \quad (3.15)$$

From the diagonalization procedure of gauge boson mass matrix (here r_s is the ratio of Higgs singlet to doublet vacuum expectation values), the following relations are obtained :-

$$M_Z^2 \cos^2 \alpha_z + M_{Z'}^2 \sin^2 \alpha_z = \frac{M_W^2}{\cos^2 \theta_z} \quad (3.16)$$

$$M_{Z'}^2 \cos^2 \alpha_z + M_Z^2 \sin^2 \alpha_z = M_W^2 \tan^2 \theta_x \times (r_s^2 + x_\phi'^2) \quad (3.17)$$

$$(M_{Z'}^2 - M_Z^2) \sin 2\alpha_z = \frac{2x_\phi' \tan \theta_x M_W^2}{\cos \theta_w} \quad (3.18)$$

From equations 3.16, 3.17 and 3.18, internal parameters dm_1 , dm_2 , xpp_2 and xpp are defined at the line numbers **264**, **273**, **283** and **292** respectively. These are given by :-

$$dm_1 = 2(M_{Z'}^2 - M_Z^2) \times sa_z \times \sqrt{1 - sa_z^2} \quad (3.19)$$

$$dm_2 = M_{Z'}^2 \cos^2 \alpha_z + M_Z^2 \sin^2 \alpha_z \quad (3.20)$$

$$xpp_2 = \frac{r_s^2 \cos^2 \theta_w \times dm_1^2}{4M_W^2 \times dm_2 - \cos^2 \theta_w \times dm_1^2} \quad (3.21)$$

$$xpp = \sqrt{xpp_2} \quad (3.22)$$

```

234 (*=====gw Modified=====*)
235 gw == {
236   ParameterType -> Internal,
237   Definitions   -> {gw->(2*MW)/vev},
238   InteractionOrder -> {QED,1},
239   TeX           -> Subscript[g,w],
240   Description    -> "Weak coupling constant at the Z pole"
241 },
242
243
244 (*=====g1 Modified=====*)
245 g1 == {
246   ParameterType -> Internal,
247   Definitions   -> {g1->gw*tw},
248   InteractionOrder -> {QED,1},
249   TeX           -> Subscript[g,1],
250   Description    -> "U(1)Y coupling constant at the Z pole"
251 },
252
253 (*=====vevs added=====*)
254 vevs == {
255   ParameterType -> Internal,
256   Value         -> rs*vev,
257   Description    -> "Singlet Scalar vacuum expectation value"
258 },
259
260 (*=====dm1 Added=====*)
261 dm1 == {
262   ParameterType -> Internal,
263   Value         -> (MZp*MZp-MZ*MZ)*2*saz*(Sqrt[1-saz*saz]),
264   TeX           -> Subscript[dm,1],
265   Description    -> "Extra Paramter 2 for Diagonalization"
266 },
267
268 (*=====dm2 Added=====*)
269 dm2 == {
270   ParameterType -> Internal,
271   Value         -> MZp*MZp*(1-saz*saz)+MZ*MZ*saz*saz,
272   TeX           -> Subscript[dm,2],
273   Description    -> "Extra Paramter 3 for Diagonalization"
274 },
275
276 (*=====xpp2 added=====*)
277 xpp2 == {
278   ParameterType -> Internal,
279   Definitions   -> {xpp2->(rs*rs*cw*cw*dm1*dm1)/(4*MW*MW*dm2-cw*cw*dm1*dm1)},
280   TeX           -> Subscript[xpp,2],
281   Description    -> "Square of Modified U(1)X Hypercharge of Higgs Doublet"
282 },
283
284 (*=====xpp added=====*)
285 xpp == {
286   ParameterType -> Internal,
287   Definitions   -> {xpp->Sqrt[xpp2]},
288   TeX           -> xpp,
289   Description    -> "Modified U(1)X Hypercharge of Higgs Doublet"
290 },
291
292

```

Figure 3.7: Parameter file describing gw, g1, vevs, dm1, dm2, xpp2 and xpp (Screenshot 5)

```

299 (*=====gxp added=====*)
300 gxp == {
301   ParameterType -> Internal,
302   Definitions   -> {gxp->gw*tx},
303   InteractionOrder -> {QED,1},
304   TeX           -> Subscript[g,xp],
305   Description    -> "Modified U(1)X coupling constant"
306 },
307
308 (*=====chi Added=====*)
309
310 chi == {
311   ParameterType -> Internal,
312   Value        -> ArcTan[(-xpp*gxp)/g1],
313   TeX          -> chi,
314   Description   -> "chi Parameter"
315 },
316
317 (*=====sc added=====*)
318
319 sc == {
320   ParameterType -> Internal,
321   Definitions   -> {sc->Sin[chi]},
322   TeX          -> sc,
323   Description   -> "Sin of Chi Parameter"
324 },
325
326 (*=====tc added=====*)
327
328 tc == {
329   ParameterType -> Internal,
330   Definitions   -> {tc->Tan[chi]},
331   TeX          -> tc,
332   Description   -> "Tan of Chi Parameter"
333 },
334
335 (*=====gx added=====*)
336
337 gx == {
338   ParameterType -> Internal,
339   Definitions   -> {gx->gxp*(1-sc*sc)},
340   InteractionOrder -> {QED,1},
341   TeX          -> Subscript[g,x],
342   Description   -> "U(1)X coupling constant"
343 },
344

```

Figure 3.8: Parameter file describing tx, gxp, chi, sc, tc and gx (Screenshot 6)

We know , the relation between modified coupling constant of $U(1)_X$ group (g'_x) and weak coupling constant (g_w) as :

$$\tan \theta_x = \frac{g'_x}{g_w} \quad (3.23)$$

So, g'_x is defined in parameter file by g_{xp} at line number **302** using equation 60 as :

$$g_{xp} = g_w \times tx \quad (3.24)$$

We know,

$$g'_x = \frac{g_x}{\cos \chi} \quad (3.25)$$

$$x'_\phi = x_\phi - \frac{g_1 \times \tan \chi}{g'_x} \quad (3.26)$$

Using 3.26 and 3.27, χ is expressed as :

$$\chi = \arctan \frac{(x'_\phi - x_\phi)g'_x}{g_1} \quad (3.27)$$

It is defined at the line number **312**. The internal parameters sc , tc and g_x are defined at the line numbers **321**, **330** and **340** respectively as :-

$$sc = \sin \chi \quad (3.28)$$

$$tc = \tan \chi \quad (3.29)$$

$$g_x = g'_x \times \sqrt{1 - sc^2} \quad (3.30)$$

$'g'_x$ is the coupling constant of $U(1)_X$ group (At the line number 41 in ZPrime.fr). Using the equations 2.36, 2.37 and 2.38, ls , μ_S , l_{HS} , lh , μ_H are defined at the line numbers **351**, **361**, **370**, **381** and **390** respectively. $'\mu'_S$ and $'\mu'_H$ are the coefficients of quadratic piece of the scalar singlet potential and scalar doublet potential (scalar sector) respectively. $'ls'$ and $'lh'$ are the coefficients of quartic piece of the scalar singlet potential and scalar doublet potential respectively. $'l''_{HS}$ is the coefficient of mixing term between scalar singlet and scalar

doublet fields in the scalar potential part of the Lagrangian. These are given by :-

$$l_s = \frac{st_s^2 \times M_H^2 + (1 - st_s^2) \times M_{H'}^2}{2 \times vevs \times vev} \quad (3.31)$$

$$\mu_S = \sqrt{vevs^2 \times l_s + \frac{l_{HS} \times vevs^2}{2}} \quad (3.32)$$

$$l_{HS} = \frac{\sqrt{1 - st_s^2} \times st_s \times M_H^2 \times M_{H'}^2}{vev \times vevs} \quad (3.33)$$

$$lh = \frac{(1 - st_s^2)M_H^2 + st_s^2 \times M_{H'}^2}{2 \times vev^2} \quad (3.34)$$

$$\mu_H = \sqrt{vev^2 \times lh + \frac{l_{HS} \times vev^2}{2}} \quad (3.35)$$

```

347 (*=====ls Added=====*)
348
349 ls == {
350   ParameterType -> Internal,
351   Value         -> (st*st*MH*MH+(1-st*st) MHP*MHP)/(vevs*vevs*2),
352   InteractionOrder -> {QED,2},
353   TeX           -> ls,
354   Description    -> "Singlet Scalar Quartic Coupling"
355 },
356
357 (*=====muS Added=====*)
358
359 muS == {
360   ParameterType -> Internal,
361   Value         -> Sqrt[vevs^2 ls+(lHS*vevs*vevs/2)],
362   TeX           -> Subscript[\[Mu],s],
363   Description    -> "Coefficient of the quadratic piece of the Singlet Scalar Potential"
364 },
365
366 (*=====lHS Added=====*)
367
368 lHS == {
369   ParameterType -> Internal,
370   Value         -> (Sqrt[1-(st^2)]*st*(MH^2-MHP^2))/(vev*vevs),
371   InteractionOrder -> {QED, 2},
372   TeX           -> Subscript[l,HS],
373   Description    -> "Coefficient of the mixing of the Singlet Scalar and Higgs field"
374 },
375
376 (*=====lh modified =====*)
377
378 lh == {
379   ParameterType -> Internal,
380   Value         -> ((1-st^2)*MH*MH + st*st*MHP*MHP)/(2*vev*vev),
381   InteractionOrder -> {QED, 2},
382   TeX           -> Higgs quartic coupling"
383 },
384
385 (*===== muH modified=====*)
386
387 muH == {
388   ParameterType -> Internal,
389   Value         -> Sqrt[vev*vev*lh +(lHS*vev*vev)/2],
390   TeX           -> Subscript[\[Mu],H],
391   Description    -> "Coefficient of the quadratic piece of the Higgs potential"
392 },
393
394
395 yl == {
396   ParameterType -> Internal,
397   Indices       -> {Index[Generation], Index[Generation]},
398   Definitions    -> {yl[i ?NumericQ, j ?NumericQ] :> 0 /; (i != j)},
399   Value         -> {yl[1,1] -> Sqrt[2] yme / vev, yl[2,2] -> Sqrt[2] ymm / vev, yl[3,3] -> Sqrt[2] ymtau / vev},
400   InteractionOrder -> {QED, 1},
401   ParameterName  -> {yl[1,1] -> ye, yl[2,2] -> ym, yl[3,3] -> ytau},
402   TeX           -> Superscript[y, l],
403   Description    -> "Lepton Yukawa couplings"
404 },

```

Figure 3.9: Parameter file describing ls, muS, lHS, lh, muH and yl (Screen-shot 7)

The internal parameters y_l , y_u , y_d and V_{CKM} are defined at the line numbers **399**, **410**, **421** and **434** respectively. The parameters y_l , y_u , y_d are the lepton, up-type quark and down-type quark Yukawa couplings respectively. V'_{CKM} is the CKM matrix parameter in our model.

```

406  yu == {
407      ParameterType -> Internal,
408      Indices       -> {Index[Generation], Index[Generation]},
409      Definitions    -> {yu[i ?NumericQ, j ?NumericQ] :> 0 /; (i != j)},
410      Value         -> {yu[1,1] -> Sqrt[2] ymup/vev, yu[2,2] -> Sqrt[2] ymc/vev, yu[3,3] -> Sqrt[2] ymt/vev},
411      InteractionOrder -> {QED, 1},
412      ParameterName  -> {yu[1,1] -> yup, yu[2,2] -> yc, yu[3,3] -> yt},
413      TeX           -> Superscript[y, u],
414      Description    -> "Up-type Yukawa couplings"
415  },
416
417  yd == {
418      ParameterType -> Internal,
419      Indices       -> {Index[Generation], Index[Generation]},
420      Definitions    -> {yd[i ?NumericQ, j ?NumericQ] :> 0 /; (i != j)},
421      Value         -> {yd[1,1] -> Sqrt[2] ymdo/vev, yd[2,2] -> Sqrt[2] yms/vev, yd[3,3] -> Sqrt[2] ymb/vev},
422      InteractionOrder -> {QED, 1},
423      ParameterName  -> {yd[1,1] -> ydo, yd[2,2] -> ys, yd[3,3] -> yb},
424      TeX           -> Superscript[y, d],
425      Description    -> "Down-type Yukawa couplings"
426  },
427
428  (* N. B. : only Cabibbo mixing! *)
429
430  CKM == {
431      ParameterType -> Internal,
432      Indices       -> {Index[Generation], Index[Generation]},
433      Unitary       -> True,
434      Value         -> {CKM[1,1] -> Cos[cabi], CKM[1,2] -> Sin[cabi], CKM[1,3] -> 0,
435                      CKM[2,1] -> -Sin[cabi], CKM[2,2] -> Cos[cabi], CKM[2,3] -> 0,
436                      CKM[3,1] -> 0, CKM[3,2] -> 0, CKM[3,3] -> 1},
437      TeX           -> Superscript[V, CKM],
438      Description    -> "CKM-Matrix"
439  }
440 }
441 };

```

Figure 3.10: Parameter file describing yu, yd and CKM (Screenshot 8)

4. Particle Classes

The fields used in the Lagrangian are described in model file by particle classes. The particle classes are labeled according to the spins of the particles.(E.g. S-Scalar Field ,F-Dirac Field,V-Vector Field). The options used in particle class are like Class Name, SelfConjugate, Mass, PDG, Width etc. This all information is included in file **particleszp.fr** which is as following :-

The physical vector fields (Gauge Bosons) A, Z, W, G and Zp are defined at the line numbers **7**, **19**, **31**, **45** and **61** respectively.

```

1  (* **** Particle classes **** *)
2
3  M$ClassesDescription = {
4
5  (* Gauge bosons: physical vector fields *)
6
7  V[1] == {
8    ClassName      -> A,
9    SelfConjugate   -> True,
10   Mass           -> 0,
11   Width          -> 0,
12   ParticleName    -> "a",
13   PDG            -> 22,
14   PropagatorLabel -> "a",
15   PropagatorType  -> W,
16   PropagatorArrow -> None,
17   FullName        -> "Photon"
18 },
19 V[2] == {
20   ClassName      -> Z,
21   SelfConjugate   -> True,
22   Mass           -> {MZ, 91.1876},
23   Width          -> {WZ, 2.4952},
24   ParticleName    -> "Z",
25   PDG            -> 23,
26   PropagatorLabel -> "Z",
27   PropagatorType  -> Sine,
28   PropagatorArrow -> None,
29   FullName        -> "Z"
30 },
31 V[3] == {
32   ClassName      -> W,
33   SelfConjugate   -> False,
34   Mass           -> {MW, Internal},
35   Width          -> {WW, 2.085},
36   ParticleName    -> "W+",
37   AntiParticleName -> "W-",
38   QuantumNumbers  -> {Q -> 1},
39   PDG            -> 24,
40   PropagatorLabel -> "W",
41   PropagatorType  -> Sine,
42   PropagatorArrow -> Forward,
43   FullName        -> "W"
44 },
45 V[4] == {
46   ClassName      -> G,
47   SelfConjugate   -> True,
48   Indices         -> {Index[Gluon]},
49   Mass           -> 0,
50   Width          -> 0,
51   ParticleName    -> "g",
52   PDG            -> 21,
53   PropagatorLabel -> "G",
54   PropagatorType  -> C,
55   PropagatorArrow -> None,
56   FullName        -> "G"
57 },
58

```

Figure 3.11: Particle file describing physical vector fields A, Z, W and G (Screenshot 1)

Using equations 2.2, 2.21 , 2.22 and 2.23, unphysical vector fields B , B_p , X , X_p and W_i (W_1 , W_2 and W_3) are defined at the line numbers **85**, **94**, **103**, **112** and **123** respectively. Fields W_1 and W_2 combine to give physical fields W^+ and W^- (As in the Standard Model)

```

59 (* =====Zp Added===== *)
60
61 V[5] == {
62   ClassName      -> Zp,
63   SelfConjugate  -> True,
64   Mass           -> {MZp,500},
65   Width          -> {WZp,10},
66   ParticleName   -> "Zp",
67   PDG            -> 32,
68   PropagatorLabel -> "Zp",
69   PropagatorType  -> Sine,
70   PropagatorArrow -> None,
71   FullName       -> "Zp"
72 },
73
74 (* Ghost fields for physical vector fields not added *)
75
76 (* Gauge bosons: unphysical vector fields *)
77
78 (* =====New definition of B Added===== *)
79
80 V[11] == {
81   ClassName      -> B,
82   Unphysical      -> True,
83   SelfConjugate   -> True,
84   Definitions     -> { B[mu_] -> Bp[mu]-tc Xp[mu]}
85 },
86
87 (* =====Bp Added===== *)
88
89 V[12] == {
90   ClassName      -> Bp,
91   Unphysical      -> True,
92   SelfConjugate   -> True,
93   Definitions     -> { Bp[mu_] -> cw A[mu] - sw Sqrt[1-saz*saz] Z[mu] + sw*saz Zp[mu]}
94 },
95
96 (* =====X Added===== *)
97
98 V[13] == {
99   ClassName      -> X,
100  Unphysical      -> True,
101  SelfConjugate   -> True,
102  Definitions     -> { X[mu_] -> (1/Sqrt[1-sc*sc]) Xp[mu]}
103 },
104
105 (* =====Xp Added===== *)
106
107 V[14] == {
108   ClassName      -> Xp,
109   Unphysical      -> True,
110   SelfConjugate   -> True,
111   Definitions     -> { Xp[mu_] -> saz Z[mu] + Sqrt[1-saz*saz] Zp[mu]}
112 },
113
114

```

Figure 3.12: Particle file describing fields Z_p , B , B_p , X and X_p (Screenshot 2)

Fields ν_l and l are defined at the line numbers **132** and **149** respectively. Physical fields ' ν_l ' and ' l ' are the neutrino and lepton fields (for three generations) respectively.

```

115 (* =====New definition of W3 Added===== *)
116
117 V[15] == {
118   ClassName      -> Wi,
119   Unphysical     -> True,
120   SelfConjugate  -> True,
121   Indices        -> {Index[SU2W]},
122   FlavorIndex    -> SU2W,
123   Definitions    -> { Wi[mu_,1] -> (Wbar[mu]+W[mu])/Sqrt[2], Wi[mu_,2] -> (Wbar[mu]-W[mu])/(I*Sqrt[2]),
124     Wi[mu_,3] -> sw A[mu]+cw Sqrt[1-saz*saz] Z[mu]-cw*saz Zp[mu]}
125 },
126
127 (* Ghost fields for unphysical vector fields not added *)
128
129
130 (* Fermions: physical fields *)
131
132 F[1] == {
133   ClassName      ->  $\nu_l$ ,
134   ClassMembers   -> {ve,vm,vt},
135   Indices        -> {Index[Generation]},
136   FlavorIndex    -> Generation,
137   SelfConjugate  -> False,
138   Mass           -> 0,
139   Width          -> 0,
140   QuantumNumbers -> {LeptonNumber -> 1},
141   PropagatorLabel -> {"v", "ve", "vm", "vt"},
142   PropagatorType  -> S,
143   PropagatorArrow -> Forward,
144   PDG             -> {12,14,16},
145   ParticleName    -> {"ve", "vm", "vt"},
146   AntiParticleName -> {"ve~", "vm~", "vt~"},
147   FullName        -> {"Electron-neutrino", "Mu-neutrino", "Tau-neutrino"}
148 },
149 F[2] == {
150   ClassName      ->  $l$ ,
151   ClassMembers   -> {e, mu, ta},
152   Indices        -> {Index[Generation]},
153   FlavorIndex    -> Generation,
154   SelfConjugate  -> False,
155   Mass           -> {ML, {Me,5.11*^-4}, {MMU,0.10566}, {MTA,1.777}},
156   Width          -> 0,
157   QuantumNumbers -> {Q -> -1, LeptonNumber -> 1},
158   PropagatorLabel -> {"l", "e", "mu", "ta"},
159   PropagatorType  -> Straight,
160   PropagatorArrow -> Forward,
161   PDG             -> {11, 13, 15},
162   ParticleName    -> {"e-", "mu-", "ta-"},
163   AntiParticleName -> {"e+", "mu+", "ta+"},
164   FullName        -> {"Electron", "Muon", "Tau"}
165 },

```

Figure 3.13: Particle file describing fields W_i , ν_l and l (Screenshot 3)

Fields uq , dq , LL and lR are defined at the line numbers **166**, **183**, **285** and **215** respectively. 'LL' fields are lepton doublets defined for all the three generations. Fields in 'LL' field are operated with **ProjM[,]** to get the left handed part of the field. 'uq' and 'dq' are the physical fields corresponding to up-type quarks and down-type quarks respectively. 'lR' are the right handed lepton fields (for electron, muon and tau). The right handed neutrino is not included in our Z' model. These are given by :-

$$LL = \begin{pmatrix} vl \\ l \end{pmatrix}_L \quad (3.36)$$

Where L stands for left handed part. Fields QL , uR , dR , H and H_p are defined at the line numbers **224**, **235**, **244**, **260** and **275** respectively. QL fields are quark doublets defined for all the three generations. 'uq' and 'dq' in 'QL' field are operated with **ProjM[,]** operator to get the left handed part of the field. For getting the right handed part of the field, we need **ProjP[,]** operator. 'uR' and 'dR' are right handed part of the fields 'uq' and 'dq' respectively (for three generations). These are given by :-

$$QL = \begin{pmatrix} uq \\ CKM \times dq \end{pmatrix}_L \quad (3.37)$$

Where L stands for left handed part of the field.

```

166 F[3] == {
167   ClassName      -> uq,
168   ClassMembers   -> {u, c, t},
169   Indices        -> {Index[Generation], Index[Colour]},
170   FlavorIndex     -> Generation,
171   SelfConjugate   -> False,
172   Mass           -> {Mu, {MU, 2.55*^-3}, {MC, 1.27}, {MT, 172}},
173   Width          -> {0, 0, {WT, 1.50833649}},
174   QuantumNumbers -> {Q -> 2/3},
175   PropagatorLabel -> {"uq", "u", "c", "t"},
176   PropagatorType  -> Straight,
177   PropagatorArrow -> Forward,
178   PDG            -> {2, 4, 6},
179   ParticleName    -> {"u", "c", "t"},
180   AntiParticleName -> {"u~", "c~", "t~"},
181   FullName       -> {"u-quark", "c-quark", "t-quark"}
182 },
183 F[4] == {
184   ClassName      -> dq,
185   ClassMembers   -> {d, s, b},
186   Indices        -> {Index[Generation], Index[Colour]},
187   FlavorIndex     -> Generation,
188   SelfConjugate   -> False,
189   Mass           -> {Md, {MD, 5.04*^-3}, {MS, 0.101}, {MB, 4.7}},
190   Width          -> 0,
191   QuantumNumbers -> {Q -> -1/3},
192   PropagatorLabel -> {"dq", "d", "s", "b"},
193   PropagatorType  -> Straight,
194   PropagatorArrow -> Forward,
195   PDG            -> {1, 3, 5},
196   ParticleName    -> {"d", "s", "b"},
197   AntiParticleName -> {"d~", "s~", "b~"},
198   FullName       -> {"d-quark", "s-quark", "b-quark"}
199 },
200
201 (* Fermions: unphysical fields *)
202
203
204
205 F[11] == {
206   ClassName      -> LL,
207   Unphysical      -> True,
208   Indices        -> {Index[SU2D], Index[Generation]},
209   FlavorIndex     -> SU2D,
210   SelfConjugate   -> False,
211   QuantumNumbers -> {Y -> -1/2, Yp -> -1/4},
212   Definitions     -> {LL[sp1_, 1, ff_] := Module[{sp2}, ProjM[sp1, sp2] vL[sp2, ff]], LL[sp1_, 2, ff_]
213     := Module[{sp2}, ProjM[sp1, sp2] l[sp2, ff]} }
214 },
215 F[12] == {
216   ClassName      -> lR,
217   Unphysical      -> True,
218   Indices        -> {Index[Generation]},
219   FlavorIndex     -> Generation,
220   SelfConjugate   -> False,
221   QuantumNumbers -> {Y -> -1, Yp -> -1/4},
222   Definitions     -> {lR[sp1_, ff_] := Module[{sp2}, ProjP[sp1, sp2] l[sp2, ff]} }
223 },

```

Figure 3.14: Particle file describing fields uq, dq, LL and lR (Screenshot 4)

Fields H and Hp are the physical scalar fields of Z' Model.

```

224 F[13] == {
225   ClassName      -> QL,
226   Unphysical     -> True,
227   Indices        -> {Index[SU2D], Index[Generation], Index[Colour]},
228   FlavorIndex    -> SU2D,
229   SelfConjugate  -> False,
230   QuantumNumbers -> {Y -> 1/6 , Yp -> 1/12},
231   Definitions    -> {
232     QL[sp1,1,ff,cc_] := Module[{sp2}, ProjM[sp1,sp2] uq[sp2,ff,cc]],
233     QL[sp1,2,ff,cc_] := Module[{sp2,ff2}, CKM[ff,ff2] ProjM[sp1,sp2] dq[sp2,ff2,cc]] }
234 },
235 F[14] == {
236   ClassName      -> uR,
237   Unphysical     -> True,
238   Indices        -> {Index[Generation], Index[Colour]},
239   FlavorIndex    -> Generation,
240   SelfConjugate  -> False,
241   QuantumNumbers -> {Y -> 2/3 , Yp -> 1/12},
242   Definitions    -> { uR[sp1,ff,cc_] := Module[{sp2}, ProjP[sp1,sp2] uq[sp2,ff,cc]] }
243 },
244 F[15] == {
245   ClassName      -> dR,
246   Unphysical     -> True,
247   Indices        -> {Index[Generation], Index[Colour]},
248   FlavorIndex    -> Generation,
249   SelfConjugate  -> False,
250   QuantumNumbers -> {Y -> -1/3 , Yp -> 1/12},
251   Definitions    -> { dR[sp1,ff,cc_] := Module[{sp2}, ProjP[sp1,sp2] dq[sp2,ff,cc]] }
252 },
253 (* Scalars *)
254 (* Higgs: Physical scalars *)
255 (* =====H Added===== *)
256 S[1] == {
257   ClassName      -> H,
258   SelfConjugate  -> True,
259   Mass           -> {MH,125},
260   Width          -> {WH,0.00407},
261   PropagatorLabel -> "H",
262   PropagatorType  -> D,
263   PropagatorArrow -> None,
264   PDG            -> 25,
265   ParticleName   -> "H",
266   FullName       -> "H"
267 },
268 (* =====Hp Added===== *)
269 S[2] == {
270   ClassName      -> Hp,
271   SelfConjugate  -> True,
272   Mass           -> {MHP,455},
273   Width          -> {WHP,0.00407},
274   PropagatorLabel -> "Hp",
275   PropagatorType  -> D,
276   PropagatorArrow -> None,
277   PDG            -> 1111,
278   ParticleName   -> "Hp",
279   FullName       -> "Hp"
280 },
281

```

Figure 3.15: Particle file describing fields QL, uR, dR, H and Hp (Screen-shot 5)

Fields hd , $G0$, GP and GZp are defined at the line numbers **298**, **303**, **317** and **336** respectively. $G0$, GP and GZp are the Goldstone bosons corresponding to Z , W and Z' bosons respectively. We know that,

$$\begin{pmatrix} hd \\ ss \end{pmatrix} = \begin{pmatrix} \cos \theta_s & -\sin \theta_s \\ \sin \theta_s & \cos \theta_s \end{pmatrix} \begin{pmatrix} H \\ Hp \end{pmatrix} \quad (3.38)$$

```

290 (* Higgs: unphysical scalars *)
291
292 (* =====New definition of hd Added===== *)
293
294 S[3] == {
295   ClassName      -> hd,
296   Unphysical     -> True,
297   SelfConjugate  -> True,
298   Definitions    -> { hd -> Sqrt[1-(st^2)] H-st Hp }
299 },
300
301
302
303 S[4] == {
304   ClassName      -> G0,
305   SelfConjugate  -> True,
306   Goldstone      -> Z,
307   Mass           -> {MZ, 91.1876},
308   Width          -> {WZ, 2.4952},
309   PropagatorLabel -> "Go",
310   PropagatorType  -> D,
311   PropagatorArrow -> None,
312   PDG            -> 250,
313   ParticleName    -> "G0",
314   FullName       -> "G0"
315 },
316
317 S[5] == {
318   ClassName      -> GP,
319   SelfConjugate  -> False,
320   Goldstone      -> W,
321   Mass           -> {MW, Internal},
322   QuantumNumbers -> {Q -> 1},
323   Width          -> {WW, 2.085},
324   PropagatorLabel -> "GP",
325   PropagatorType  -> D,
326   PropagatorArrow -> None,
327   PDG            -> 251,
328   ParticleName    -> "G+",
329   AntiParticleName -> "G-",
330   FullName       -> "GP"
331 },
332
333
334 (* =====GZp Added===== *)
335
336 S[6] == {
337   ClassName      -> GZp,
338   SelfConjugate  -> False,
339   Goldstone      -> Zp,
340   Mass           -> {MZp, 500},
341   QuantumNumbers -> {Q -> 0},
342   Width          -> {WZp, 10},
343   PropagatorLabel -> "GZp",
344   PropagatorType  -> D,
345   PropagatorArrow -> None,
346   PDG            -> 2222,
347   ParticleName    -> "GZp+",
348   AntiParticleName -> "GZp-",
349   FullName       -> "GZp"
350 },
351

```

Figure 3.16: Particle file describing fields hd , $G0$, GP and GZp (Screenshot 6)

Fields ss, Phi and Su are defined at the line numbers **359**, **372** and **383** respectively. 'hd' and 'ss' are the unphysical scalar fields and are defined using the equations 3.38. Fields Phi and Su are the scalar doublet and scalar singlet fields respectively. They are defined using the equations 2.28 and 2.29.

```

353 (* =====ss Added===== *)
354
355 S[7] == {
356   ClassName    -> ss,
357   Unphysical   -> True,
358   SelfConjugate -> True,
359   Definitions  -> { ss -> st*H+Sqrt[1-(st^2)] Hp }
360 },
361
362 (*===== Higgs doublet and Singlets=====*)
363
364 S[11] == {
365   ClassName    -> Phi,
366   Unphysical   -> True,
367   Indices      -> {Index[SU2D]},
368   FlavorIndex  -> SU2D,
369   SelfConjugate -> False,
370   QuantumNumbers -> {Y -> 1/2 , Yp -> 0},
371   Definitions  -> { Phi[1] -> -I GP, Phi[2] -> (vev+hd+I G0)/Sqrt[2] }
372 },
373
374
375 (* =====Su Added=====*)
376
377 S[12] == {
378   ClassName    -> Su,
379   Unphysical   -> True,
380   SelfConjugate -> False,
381   QuantumNumbers -> {Y -> 0 , Yp -> 1/2},
382   Definitions  -> { Su -> (vevs+ss+I GZp)/Sqrt[2]}
383 },
384
385 };
386

```

Figure 3.17: Particle file describing fields ss, Phi and Su (Screenshot 7)

3.2 Lagrangian

The Lagrangian in the Wolfram Language (in mathematica) for Z' is shown in the figure **3.18**. The Lagrangian has four parts namely LYukawa, LFermions, LGauge and LScalar. LYukawa and LFermions are the same as that of the Standard Model. The part of the Lagrangian including ghost fields is not included in the model.

```

LYukawa := Block[{sp, ii, jj, cc, ff1, ff2, ff3, yuk, feynmangaugerules}, feynmangaugerules = If[Not[FeynmanGauge],
  {G0 | GP | GPbar -> 0}, {}];
yuk = ExpandIndices[-yd[ff2, ff3] * CKM[ff1, ff2] * QLbar[sp, ii, ff1, cc].dR[sp, ff3, cc] * Phi[ii] -
  yl[ff1, ff3] * LLbar[sp, ii, ff1].lR[sp, ff3] * Phi[ii] - yu[ff1, ff2] * QLbar[sp, ii, ff1, cc].uR[sp, ff2, cc] *
  Phibar[jj] * Eps[ii, jj], FlavorExpand -> SU2D];
yuk = yuk /. {CKM[a_, b_] Conjugate[CKM[a_, c_]] -> IndexDelta[b, c], CKM[b_, a_] Conjugate[CKM[c_, a_]] -> IndexDelta[b, c]};
yuk + HC[yuk] /. feynmangaugerules
];

LGauge := Block[{mu, nu, ii, aa}, ExpandIndices[-1/4 FS[B, mu, nu] * FS[B, mu, nu] - 1/4 FS[X, mu, nu] * FS[X, mu, nu] -
  sc/2 FS[B, mu, nu] * FS[X, mu, nu] - 1/4 FS[Wi, mu, nu, ii] * FS[Wi, mu, nu, ii] - 1/4 FS[G, mu, nu, aa] *
  FS[G, mu, nu, aa], FlavorExpand -> SU2W]];

LFermions := Block[{mu}, ExpandIndices[I * (QLbar.Ga[mu].DC[QL, mu] + LLbar.Ga[mu].DC[LL, mu] + uRbar.Ga[mu].DC[uR, mu] +
  dRbar.Ga[mu].DC[dR, mu] + lRbar.Ga[mu].DC[lR, mu]), FlavorExpand -> {SU2W, SU2D}] /. {CKM[a_, b_] *
  Conjugate[CKM[a_, c_]] -> IndexDelta[b, c], CKM[b_, a_] Conjugate[CKM[c_, a_]] -> IndexDelta[b, c]};

LScalar := Block[{ii, mu, feynmangaugerules}, feynmangaugerules = If[Not[FeynmanGauge], {G0 | GP | GPbar -> 0}, {}];
ExpandIndices[DC[Phibar[ii], mu] * DC[Phi[ii], mu] + muH^2 Phibar[ii] * Phi[ii] - lH Phibar[ii] *
  Phi[ii] * Phibar[jj] * Phi[jj] + DC[Subar, mu] * DC[Su, mu] + muS^2 Subar.Su - ls (Subar.Su) * (Subar.Su) -
  lHS Phibar[ii] * Phi[ii] Subar.Su, FlavorExpand -> {SU2D, SU2W}] /. feynmangaugerules];

LZPM := LGauge + LFermions + LScalar + LYukawa;

```

Figure 3.18: Lagrangian

3.3 Running The FeynRules

The first thing that must be done when using FeynRules is to load the package into a Mathematica session. This should be done before any of the model file is loaded in the kernel. In order to load FeynRules, the user must first specify the directory where it is stored and then load it. After the FeynRules package has been loaded, the model can be loaded using the command **LoadModel[]**.

The Lagrangian is given during the mathematica session. Some points for Lagrangian are :-

1. FeynRules provides functions that can be used while constructing Lagrangians like $DC[\phi, \mu]$ for covariant derivative, $Ga[\mu]$ for Gamma(Dirac) matrix etc.
2. Once the Lagrangian is implemented, several checks are performed by means of the functions like CheckHermiticity[L], GetInteractionTerms[L], GetMassTerms[L] etc.

After the model description is created and the Lagrangian constructed,

it can be loaded into FeynRules and the Feynman rules are obtained. The command used is **FeynmanRules[L]**.

UFO Interface

After issuing this command, UFO files needed for MadGraph5 are created using command **WriteUFO[L]**. These files are basically python files(**.py**) files which contain everything about model like parameters, vertices, particles etc. With the UFO files, there is a **logfile** for our model where we can check if there are errors in our model like electric charges are defined correctly or not, lepton number conservation etc.

Chapter 4

MadGraph

MadGraph[6] is a software package which allows physicists to study the properties of particles and their interactions in a simulated environment that can be used to make predictions for experimental results. MadGraph uses Monte Carlo techniques to generate events, which means that it uses random numbers to simulate the behavior of particles in a collision. The simulated events can be compared to experimental data to test theoretical models and make predictions for future experiments. It can be used for finding scattering cross-section and generating events.

4.1 Installation

Here , I have used MadGraph5aMC@NLO. MadGraph5aMC@NLO is the new version of both MadGraph5 and aMC@NLO that unifies the LO and NLO lines of development of automated tools within the MadGraph family. The installation procedure is shown in the figure 4.1

Website: <https://launchpad.net/mg5amcnlo>

Software Requirements

Python 3.9, gfortran compiler

Others

MadAnalysis5, lhapdf6 (These can be installed in MadGraph by **Install** command)

The various steps required for generating the events using MadGraph

```

root:~$ sudo apt install software-properties-common
root:~$ sudo apt-get install --reinstall ca-certificates
root:~$ sudo add-apt-repository ppa:deadsnakes/ppa
root:~$ sudo apt update
root:~$ sudo apt install python3.9
root:~$ sudo apt install gfortran
root:~$ python3.9 --version
root:~$ cd Desktop/Joginder/Project/Mad_Graph

tar -xvzf MG5_aMC_v3.4.1.tar.gz

root:~$ Desktop/Joginder/Project/Mad_Graph/MG5_aMC_v3_4_1/bin$ ./mg5_aMC

```

Figure 4.1: Installation

are shown in the figure 4.6. These are explained as follows :-

1. **Importing model** The UFO Files folder created in the model folder must be copied to the directory of MadGraph (models folder of bin). Then, it can be imported by issuing command : **import model modelname**.
2. **Generating a Process** A process can be generated by **generate p1 p2 > p3 p4**. There must be a space between different symbols of particles used in MadGraph. We can add different processes together. We want a process where somehow Z' boson is involved. There are different ways by which Z' particle can be produced and decayed. It can be produced by pp collision or electron-positron collision. Similarly, it can decay to dilepton pair or to W^+ and W^- particles. But we have chosen the process in equation 76. I have used the experimental data of Large Hadron Collider with the ATLAS detector[2]. LHC is a hadron collider and so have used pp collision to produce Z' . The decay width of Z' to dilepton state is maximum or is the dominant channel in Z' decay.

$$pp > zp > e + e^- \quad (4.1)$$

When we upload the UFO files of a particular model, the particle def-

initions and symbols are defined. Here p means proton, zp means Z' boson and $e^+ e^-$ means positron and electron respectively. In terms of MadGraph, two new particles are included in Z' model apart from the particles in the Standard model (namely **zp** and **hp**).

3. Feynman Diagrams .

The diagrams can be displayed by issuing the command : **display diagrams**. The Feynman diagrams for our selected process are as shown below :-

`u u~ > zp > e+ e- WEIGHTED=4`

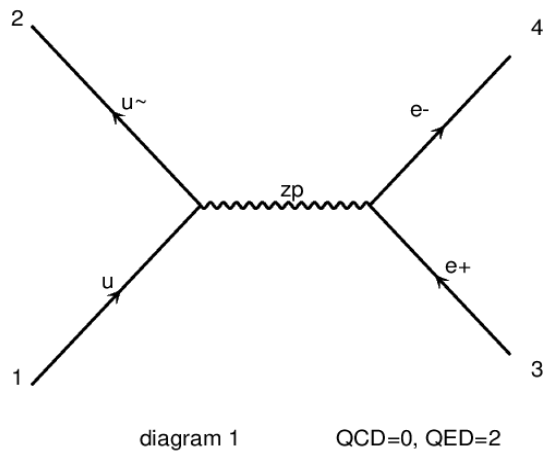


Figure 4.2: Feynman Diagram (Screenshot 1)

c c~ > zp> e+ e- WEIGHTED=4

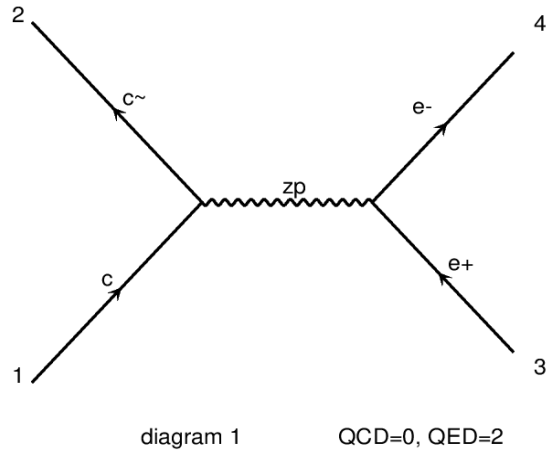


Figure 4.3: Feynman Diagram (Screenshot 2)

s s~ > zp> e+ e- WEIGHTED=4

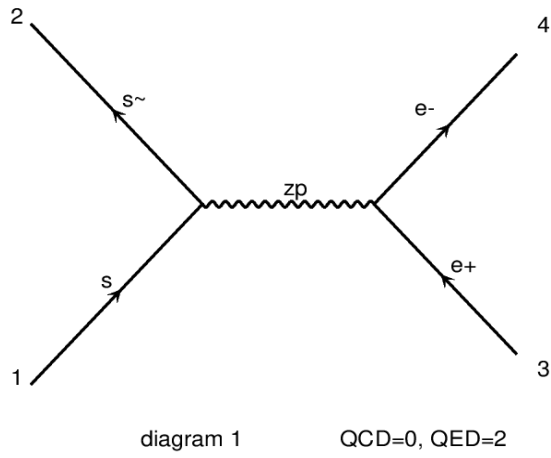


Figure 4.4: Feynman Diagram (Screenshot 3)

d d~ > zp > e+ e- WEIGHTED=4

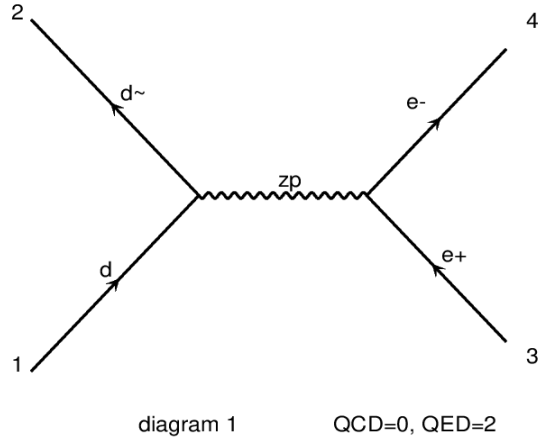


Figure 4.5: Feynman Diagram (Screenshot 4)

4. **Launching the Process** For event generation, our whole event data must be stored in a folder. This is done by **output foldername**(here output zp1). Then we have to launch process by Launch command (**launch foldername**). There are various sub folders in this folder like Cards, Events, bin etc.

```
MG5_aMC>import model Z_Prime_Model_UFO
MG5_aMC>generate p p > zp > e+ e-
MG5_aMC>display diagrams
MG5_aMC>output zp1
MG5_aMC>launch zp1
MG5_aMC>open index.html
```

Figure 4.6: Event Generation

After issuing the command **open index.html**, the results page will be displayed on the browser screen.

A run for our selected process is given by :-

Results in the Z_Prime_Model_UF01 for $p p > z p > e^+ e^-$

Available Results

Run	Collider	Banner	Cross section (pb)	Events	Data	Output	Action
run_01	$p p$ 6500.0 x 6500.0 GeV	tag_1	1.487 ± 0.00093	100000	parton madevent	LHE MA5_report_analysis1	remove run launch detector

[Main Page](#)

Figure 4.7: Sample Run

4.2 Cards Files

Cards files are the input files to the MadGraph. They include all the information regarding event generation and the chosen process. There are many types of card files. The commonly needed ones are run card, parameter card, Mad-analysis parton card, plot card etc. These files can be changed individually or during the run time. For our selected process, the following card files are changed :-

1. **Run Card** In the run card, the number of events and the energy of colliding beams are changed. There are other things also but they are already set by default like rapidity , minimum invariant mass of dilepton pair etc. An example is as following :-

100000 = nevents ! Number of unweighted events requested

6500.0 = ebeam1 ! beam 1 total energy in GeV

6500.0 = ebeam2 ! beam 2 total energy in GeV

2. **Parameter Card** Parameter card contains information of the various parameters of the model. We can change the value of external parameters during the actual run and don't need to upload the model again and again.

3. **Mad-Analysis Parton Card** This card includes the information

regarding plot range and luminosity. An example is as following :-

```
set main.lumi = 36
```

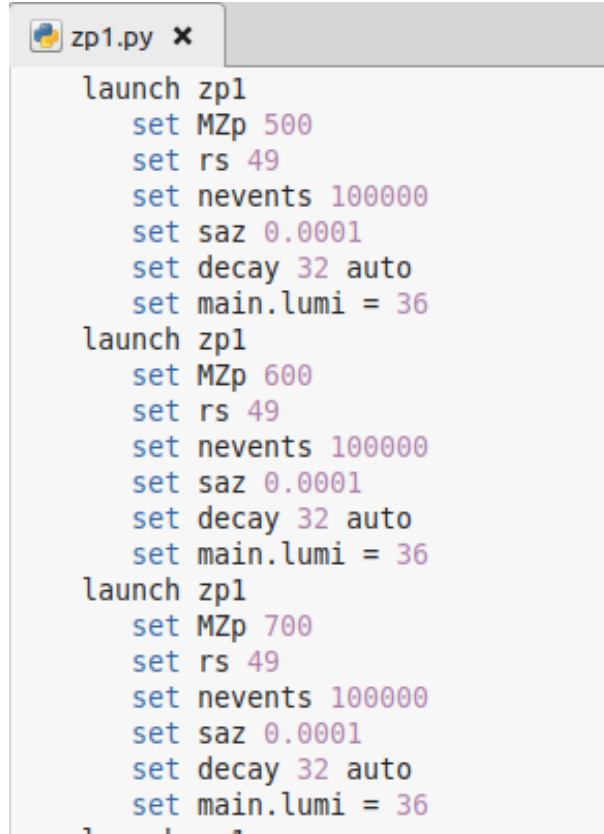
```
plot M(e-[1] e+[1]) 40 0 500 [logY ]
```

4. **Plot Card** This card includes the information of plots and their definitions in the MadGraph.

4.3 Plots

For every run of the process, we get a cross-section value in picobarns. Each run consists of a number of events which information is given in the run card. The total scattering cross-section depends on various factors like centre-of-mass energy of colliding particles, mass of a particular particle etc. Here, we are changing the mass of Z' particle (MZp name in FeynRules code). Further, decay width of Z' particle is also changed as the mass is changed. A single run is done for a particular value of $\sin(\alpha_z)$ and rs (ratio of singlet to doublet vacuum expectation values). As MadGraph takes input as python files, we can make a script file written in python to have runs queued one after other. Since mass of Z' particle is changed from 0.5 TeV to 5 TeV, part of script file code for 500 GeV to 700 GeV is as shown below :-

Instead of making a script file, we can also give range of the particular parameter in the parameter card file. We have changed mass of Z' particle (or MZp parameter in the code with PDG code of 32) from 500 GeV to 5000 GeV in our runs. It is given by :-



```
launch zp1
  set MZp 500
  set rs 49
  set nevents 100000
  set saz 0.0001
  set decay 32 auto
  set main.lumi = 36
launch zp1
  set MZp 600
  set rs 49
  set nevents 100000
  set saz 0.0001
  set decay 32 auto
  set main.lumi = 36
launch zp1
  set MZp 700
  set rs 49
  set nevents 100000
  set saz 0.0001
  set decay 32 auto
  set main.lumi = 36
```

Figure 4.8: Script File

We will be plotting mainly two types of diagram here. The first one is cross-section times branching ratio ($\sigma\mathbf{B}$ in pico barn) vs mass of Z' boson ($\mathbf{MZp}$ in TeV) diagram. This type of graph is compared with the experimental data of LHC(Large Hadron Collider) ALICE detector[2] ($\sqrt{s} = 13TeV, 36fb^{-1}$). When compared, there will be a point where both the graph will intersect. This will give a lower limit on the mass of Z' particle. When sine of Z- Z' mixing angle ($\sin(\alpha_z)$) is changed, we get a different lower limit on Z' particle mass ($\mathbf{MZp}$). So, the second type of graph is plotted between $\mathbf{MZp}$ vs $\sin(\alpha_z)$.

```
#####
## PARAM_CARD AUTOMATICALLY GENERATED BY MG5      ###
#####
#####
## INFORMATION FOR CKMBLOCK
#####
BLOCK CKMBLOCK #
  1 2.277360e-01 # cabi
#####
## INFORMATION FOR MASS
#####
BLOCK MASS #
  1 5.040000e-03 # md
  2 2.550000e-03 # mu
  3 1.010000e-01 # ms
  4 1.270000e+00 # mc
  5 4.700000e+00 # mb
  6 1.720000e+02 # mt
 11 5.110000e-04 # me
 13 1.056600e-01 # mmu
 15 1.777000e+00 # mta
 23 9.118760e+01 # mz
 25 1.250000e+02 # mh
 32 scan: [500,600,700,800,900,1000,1100,1200,1300,1400,1500,1600,1700,1800,1900,2000,2100,2200,2300,
2400,2500,2600,2700,2800,2900,3000,3100,3200,3300,3400,3500,3600,3700,3800,3900,4000,
4100,4200,4300,4400,4500,4600,4700,4800,4900,5000] # 3.500000e+03 # mzp
1111 4.550000e+02 # mhp
 12 0.000000e+00 # ve : 0.0
 14 0.000000e+00 # vm : 0.0
 16 0.000000e+00 # vt : 0.0
 21 0.000000e+00 # g : 0.0
 22 0.000000e+00 # a : 0.0
 24 7.982436e+01 # w+ : cw*mz
.....
```

Figure 4.9: Parameter Card

Chapter 5

Results

The data set for all the runs for $\sin(\alpha_z)=0.0001$ and the corresponding graph is shown in the figures 5.2 and 5.1 respectively. The experimental plot[2] in red colour is shown in figure 5.3. This graph tells us that the cross-section values have a maximum (upper limit) for a particular $M_{Z'}$. Only the values of cross-section which are lower than this maximum value are possible. So, when the model and experimental graphs are compared, a lower limit on $M_{Z'}$ is found near the intersection point. The first type of graph between $\sigma_{\mathbf{B}}$ and $\mathbf{MZp}$ is shown for $\sin(\alpha_z)=0.0001$ in the figure 5.5.

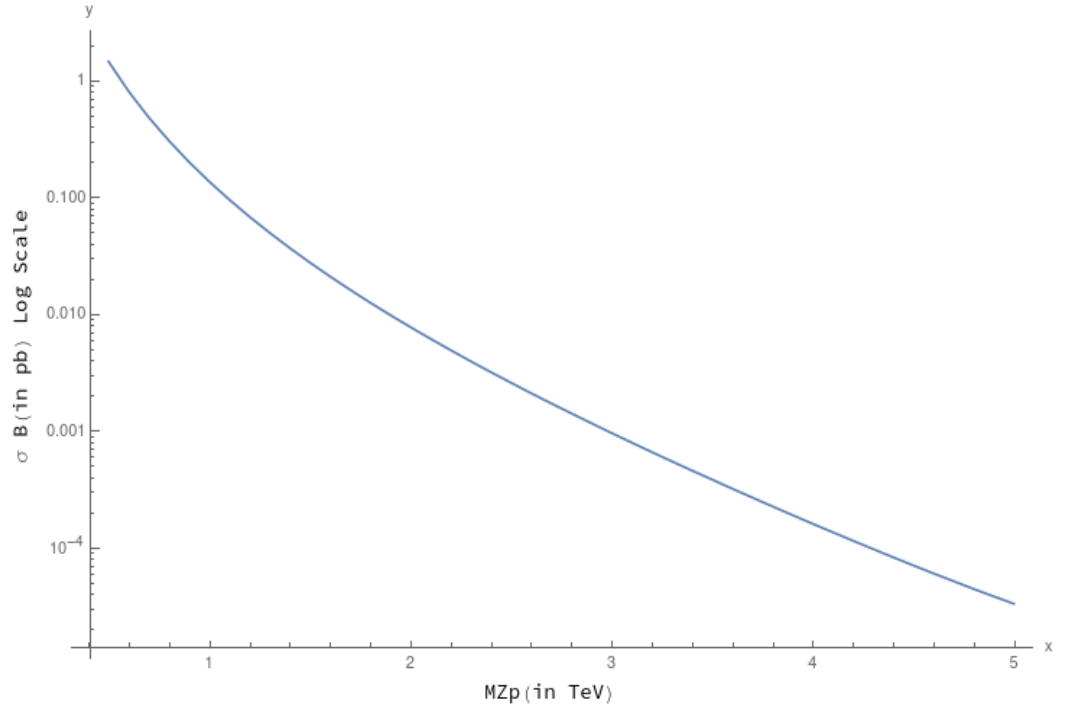


Figure 5.1: Model Data Plot

1	MZp(TeV) GIVEN Cross-Section(pb) (rs=40)	
2	0.5	1.446
3	0.6	0.8027
4	0.7	0.4801
5	0.8	0.3032
6	0.9	0.1995
7	1	0.1357
8	1.1	0.09496
9	1.2	0.06785
10	1.3	0.04951
11	1.4	0.03662
12	1.5	0.02751
13	1.6	0.0209
14	1.7	0.01607
15	1.8	0.01246
16	1.9	0.009751
17	2	0.007694
18	2.1	0.006105
19	2.2	0.004879
20	2.3	0.003917
21	2.4	0.003162
22	2.5	0.002566
23	2.6	0.002087
24	2.7	0.001708
25	2.8	0.001402
26	2.9	0.001154
27	3	0.0009522
28	3.1	0.0007878
29	3.2	0.0006535
30	3.3	0.000544
31	3.4	0.0004533
32	3.5	0.0003797
33	3.6	0.0003174
34	3.7	0.0002664
35	3.8	0.0002237
36	3.9	0.0001883
37	4	0.0001587
38	4.1	0.0001344
39	4.2	0.0001137
40	4.3	9.65E-05
41	4.4	8.20E-05
42	4.5	6.99E-05
43	4.6	5.97E-05
44	4.7	5.12E-05
45	4.8	4.39E-05
46	4.9	3.79E-05
47	5	3.28E-05

Figure 5.2: Model Data Set

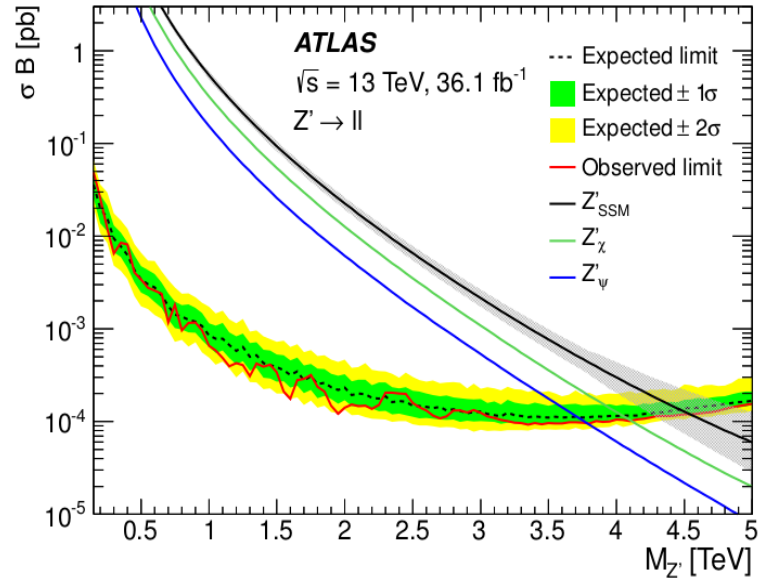


Figure 5.3: Experimental Data Plot 1

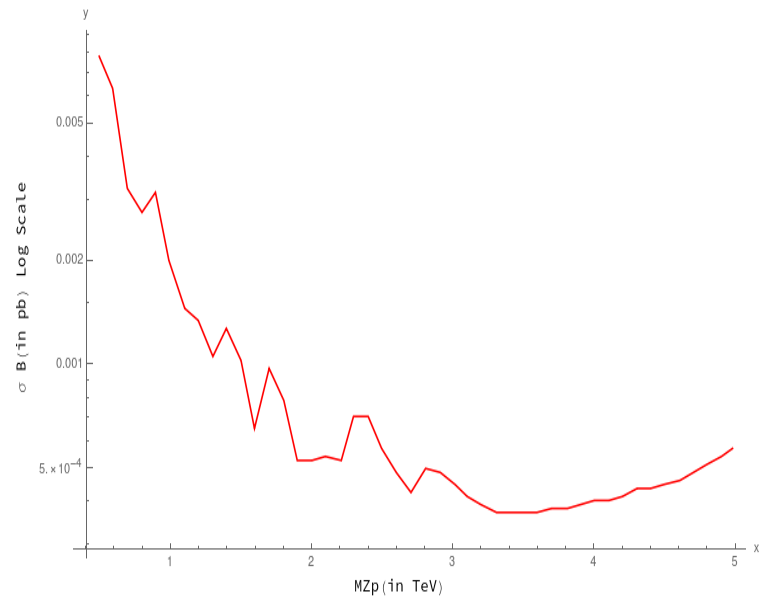


Figure 5.4: Experimental Data Plot 2

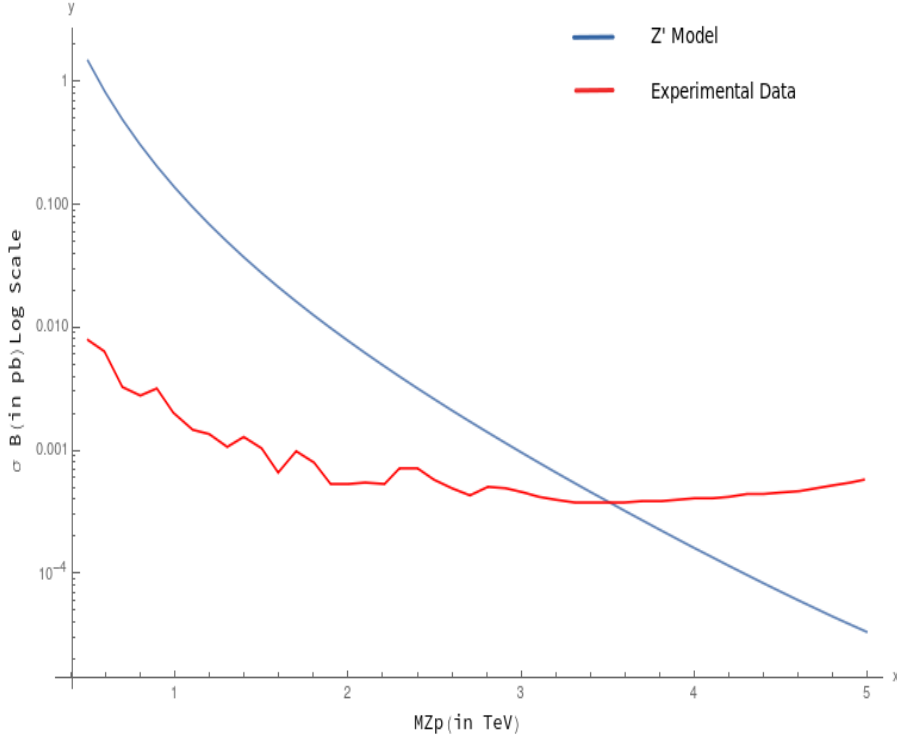


Figure 5.5: Combined Plot

The intersection point for the above graph (figure 5.5) is at $MZp=3.50332$ TeV. So, according to this, Z' particle cannot have a mass value lower than 3.50332 TeV. Mass limits on Z' are found for different values of $\sin(\alpha_z)$ parameter. In all these runs, rs and $\tan\theta_x$ values are fixed to 40 and 1 respectively. A graph can be plotted between MZp and $\sin(\alpha_z)$. The data set points and the graph are shown in the figures 5.6 and 5.7 respectively.

$\sin(\alpha_z) \times 10^4$	$MZp(\text{Found})$
-4	4.77503
-3	4.3091
-2	3.91991
-1	3.5088
0	2.93324
1	3.50332
2	3.91443
3	4.3091
4	4.786

Figure 5.6: Model Data Set 2

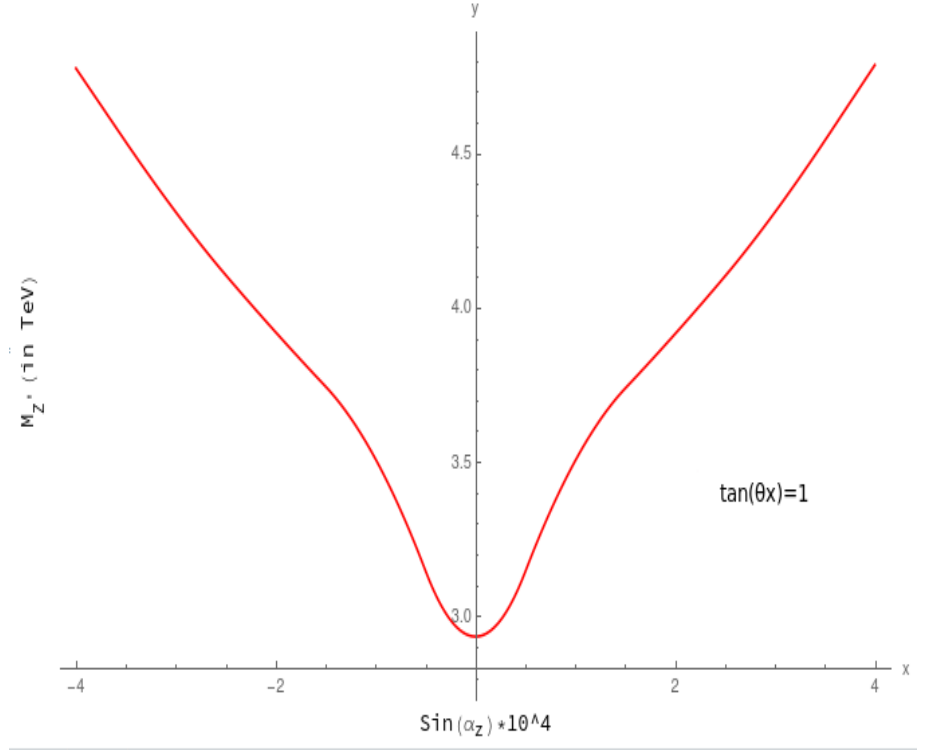


Figure 5.7: Final Plot

5.1 Future

Experimental data has been updated for greater center-of-mass energy and luminosity. We can find the new lower Z' particle mass limits for new experimental data of hadron collider and electron-positron collider.

References

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- [3] T. Bandyopadhyay, G. Bhattacharyya, D. Das, and A. Raychaudhuri, “Reappraisal of constraints on Z models from unitarity and direct searches at the LHC,” *Physical Review D*, vol. 98, no. 3, p. 035027, 2018.
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